

Calabi-Yau Fourfolds with non-trivial Three-Form Cohomology

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arXiv: 1512.04859, 1702.03217 (T. Grimm, SG)

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Motivation

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F-theory effective action on elliptically fibered Calabi-Yau fourfold Y_4 :

$\mathcal{N} = 1$ gauged supergravity in $(3 + 1)$ dimensions

$$S^{(4)} = \int_{\mathcal{M}_{3,1}} \frac{1}{2} R^{(4)} * 1 - K_{IJ}^F \mathcal{D}M^I \wedge * \mathcal{D}\bar{M}^J \\ - \frac{1}{2} \text{Re}(f)_{\Lambda\Sigma} F^\Lambda \wedge * F^\Sigma - \frac{1}{2} \text{Im}(f)_{\Lambda\Sigma} F^\Lambda \wedge F^\Sigma$$

Massless spectrum includes **complex scalars** from **non-trivial** harmonic **three-forms** of Y_4

⇒ What can we learn about their **dynamics**?

⇒ Can we construct explicit **examples**?

Outline

- Three-Forms on Calabi-Yau Fourfolds
- Calabi-Yau Fourfolds as Toric Hypersurfaces
- Metric for Three-Form Moduli of a Toric Hypersurface
- Outlook

Geometry and Topology of CY_4

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Characteristic feature:

$SU(4)$ -holonomy

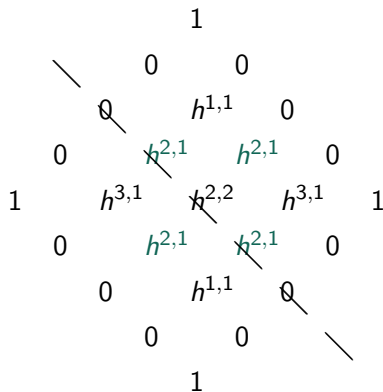
Covariantly constant Weyl-spinor:

$$\nabla\eta = 0, \quad \gamma_9\eta = \eta$$

$\Rightarrow \mathcal{N} = 2$ 3d M-theory vacua!
(no flux)

$\Rightarrow \mathcal{N} = 1$ 4d F-theory vacua!
(no flux)

Hodge diamond:



Expansion of the Three-Form

Consider scalar three-form moduli arising from $C_3 \in H^{2,1}(Y_4) \oplus H^{1,2}(Y_4)$

Kinetic term: $dC_3 \wedge *dC_3 \Rightarrow$ Hodge-star depends on moduli!

precisely: $*\psi = iJ \wedge \psi$, $\psi \in H^{2,1}(Y_4)$, $J \in H^{1,1}(Y_4)$

\Rightarrow Choose three-form basis Ψ^I depending on complex structure moduli!

$$\Psi^I(z, \bar{z}) = \frac{1}{2} \text{Re}(f)^{Im} (\alpha_m - i\bar{f}_{mk}\beta^k) \in H^{1,2}(Y_4)$$

Three-form deformations:

$$C_3 = \sum_{I=1}^{h^{2,1}} N_I \Psi^I + c.c.$$

$\Rightarrow h^{2,1}$ complex scalars N_I

\Rightarrow bosonic part of chiral multiplets

Three-Form Ansatz

Suggested by [Grimm '10]:
($l, k, m = 1, \dots, h^{2,1}$)

$$\Psi^l = \frac{1}{2} \operatorname{Re}(f)^{lm} (\alpha_m - i \bar{f}_{mk} \beta^k) \in H^{1,2}(Y_4)$$

Properties:

- α_l, β^k basis of $H^3(Y_4, \mathbb{R}) \Rightarrow$ topological
- $f_{lm}(z)$ holomorphic (three-form periods)
- Assume: $\beta_l \wedge \beta_m = 0$

Result:

$$\frac{\partial \Psi^l}{\partial z^k} = \frac{\partial \bar{\Psi}^l}{\partial z^k} \sim \Psi^m$$

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Advantage: $M_{Al}{}^k = \int \omega_A \wedge \alpha_l \wedge \beta^k$ topological intersection numbers

$$\Rightarrow \int_{Y_4} \Psi^l \wedge \bar{\Psi}^k = -\frac{1}{2} \operatorname{Re}(f)^{lm} M_{Am}{}^k v^A$$

Goal of our work:
Calculate f and M !

$$K_{new}^M \sim \operatorname{Re}(N)_l \operatorname{Re}(f)^{lm} M_{Am}{}^k v^A \operatorname{Re}(N)_k \Rightarrow \operatorname{Im}(N)_l \text{ axionic!}$$

Toric Calabi-Yau Fourfold Hypersurfaces

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Now Calabi-Yau fourfold Y_4 smooth toric hypersurface

Toric divisors D_i , codimension one submanifolds of Y_4 invariant under torus action

Non-trivial three- and two-forms from divisors D_i via Gysin-morphism of their inclusion

$$\iota_i : D_i \rightarrow Y_4 \quad [\text{Danilov;Mavlyutov}]$$

Gysin-Morphism

$$\iota_{i*} : H^n(D_i, \mathbb{C}) \xrightarrow{PD} H_{6-n}(D_i, \mathbb{C}) \xrightarrow{(\iota_i^*)^{-1}} H_{6-n}(Y_4, \mathbb{C}) \xrightarrow{PD} H^{n+2}(Y_4, \mathbb{C})$$

⇒ Respect Hodge decomposition:

$$\iota_{i*} : H^{0,0}(D_i) \longrightarrow H^{1,1}(Y_4),$$

$$\iota_{i*} : H^{1,0}(D_i) \longrightarrow H^{2,1}(Y_4).$$

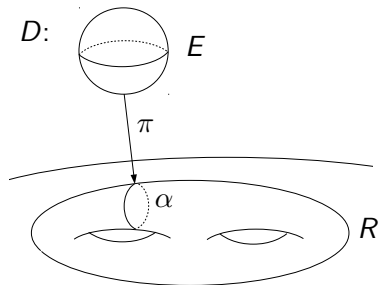
For smooth toric Calabi-Yau hypersurfaces Y_4

$$\bigoplus_i H^{0,0}(D_i) \simeq H^{1,1}(Y_4),$$

$$\bigoplus_i H^{1,0}(D_i) \simeq H^{2,1}(Y_4), \quad \bigoplus_i H^{0,1}(D_i) \simeq H^{1,2}(Y_4).$$

⇒ Construct $H^{1,0}(D_i)$ for toric divisors D_i of Y_4 !

Divisors with non-trivial One-Forms



$$H^{1,0}(D) \simeq H^{1,0}(R)$$

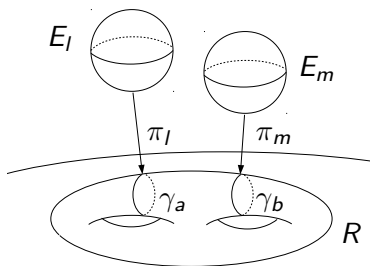
For $H^{1,0}(D) \neq 0$:

Toric divisor D fibration

Fiber: E toric surface,
($h^{i,j}(E) = 0$ for $i \neq j$)

Base: R Riemann surface

Three-forms of Calabi-Yau Fourfold Hypersurface



Construct three-forms:

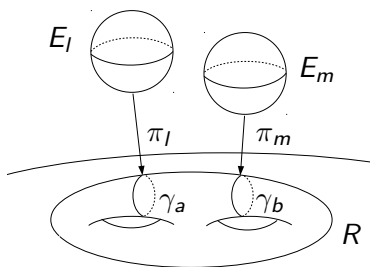
$$\psi_{\mathcal{A}} = \iota_{l*}(\gamma_a) \in H^{2,1}(Y_4)$$

with $\mathcal{A} = (l, a)$

$$\iota_l : D_l \rightarrow Y_4, \quad \gamma_a \in H^{1,0}(R)$$

\Rightarrow **Complex structure** dependence
of $\psi_{\mathcal{A}}$ determined by $H^{1,0}(R)$

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Normalized basis: $\gamma_a = \alpha_a + i f_{ab}(z) \beta^b \in H^{1,0}(R), \quad \alpha_a, \beta^b \in H^1(R, \mathbb{Z})$

Metric on $H^{1,0}(R)$:
$$-i \int_R \gamma_a \wedge \bar{\gamma}_b = 2 \cdot \text{Re}(f)_{ab} > 0, \quad f_{ab} = f_{ba}$$

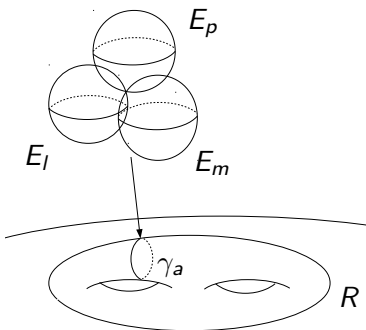
Three-Form Metric for Fourfold Hypersurface

Recall metric on $H^{2,1}(Y_4)$

$$Q_{AB} = \int_{Y_4} \psi_A \wedge * \bar{\psi}_B = -i \int_{Y_4} J \wedge \psi_A \wedge \bar{\psi}_B, \quad J \in H^{1,1}(Y_4)$$

Two- and three-forms from toric divisors via Gysin-map

\Rightarrow Triple intersection of toric divisors!



Only divisors that fiber over same R contribute!

$$D_l \cap D_m \cap D_p = M_{Imp} \cdot R$$

with $M_{Imp} = \# E_l \cap E_m \cap E_p$

(Generalized sphere tree)

Three-Form Metric for Fourfold Hypersurface

Metric on $H^{2,1}(Y_4)$

$$Q_{AB} = 2v^I M_{Imp} \operatorname{Re}(f)_{ab}$$

$\mathcal{A} = (m, a)$, $\mathcal{B} = (p, b)$

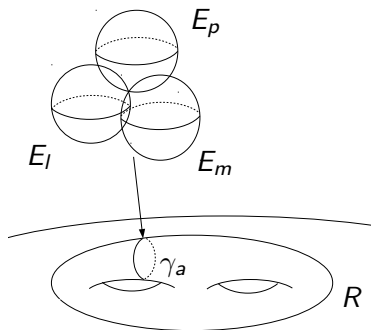
Reorder indices to find previous result

$$Q_{AB} = 2v^\Sigma M_{\Sigma\mathcal{A}}^C \operatorname{Re}(f)_{C\mathcal{B}}$$

$v^\Sigma M_{\Sigma\mathcal{A}}^C$ linear dependence on Kähler moduli v^Σ

$\operatorname{Re}(f)_{C\mathcal{B}}$ complicated dependence on complex structure moduli z^K

\Rightarrow Toric data exchanged by mirror symmetry!



Outlook

Next step:

Discuss **F-theory physics** on elliptically fibered CY_4 and their weak coupling limit on CY_3

Construct examples with three-forms such that weakly coupled IIB has

- **Wilson line moduli** on a $D7$ wrapping a four-cycle of CY_3 (axions)
- **Odd moduli**, scalars from B_2, C_2 (axions)

⇒ **Metric** on $H^{2,1}(Y_4)$ determines **decay-constants!**

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Future directions:

- Construct **bases** with three-forms (**bulk $U(1)$**)
- Generalize to **complete intersections**
- Include **G_4 -flux** and calculate **superpotential**
- ...

Thank you for your attention!

Questions?