

# Holography for Composite Inflation

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( with P. Suranyi, L.C.R. Wijewardhana )

# Cosmological Inflation:

Standard description:

- expansion driven by the potential energy of a scalar field  $\varphi$  called **inflaton**
- weakly coupled Lagrangian for the inflaton within QFT framework:

$$S = \int d^4x \sqrt{-\det g} \left[ \frac{R}{2} + \frac{1}{2} g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - V(\varphi) \right]$$

- preferably: small field models  
( $\Delta\varphi \ll M_P \Rightarrow$  EFT reliable)

BUT:

$\eta$  problem:

Recall the slow roll conditions:

$$\varepsilon = \frac{V'(\varphi)}{V(\varphi)} \ll 1 \quad , \quad \eta = \frac{V''(\varphi)}{V(\varphi)} \ll 1$$

(consistency with observations  $\Rightarrow$  slow roll inflation)

However: **Quantum corrections** drive inflaton mass ( $m_\varphi^2 = V''$ ) to cutoff of effective theory (at least Hubble scale  $H \approx \sqrt{V}$ )

$\rightarrow \Delta\eta \approx \mathcal{O}(1)$  or larger  $\Rightarrow$  inflation ends prematurely

Hence need a symmetry... (ex.: axion monodromy inflation...)

# Composite Inflation:

A possible different approach:

Inflaton - a composite state in a strongly coupled gauge theory

[F. Bezrukov, P. Channuie, J. Joergensen, F. Sannino, arXiv:1112.4054; [inflaton - glueball](#)]

→ inflaton mass dynamically fixed  $\Rightarrow$  no  $\eta$  problem!

Recently was argued that tensor-to-scalar ratio  $r$  can be large in such models [P. Channuie, K. Karwan, arXiv:1404.5879]

**Our aim:** Use Gauge/Gravity Duality (GGD) to study this class of inflationary models

## Gauge/gravity duality:

D-branes: open string BCs  $\longleftrightarrow$  SUGRA solutions

strong coupling  $\longleftrightarrow$  weak coupling

$\Rightarrow$  Can use classical supergravity to learn about strongly coupled gauge theories

### Some potential applications:

- high  $T_c$  superconductivity
- quark-gluon plasma [viscosity/entropy density]
- dynamical electroweak symmetry breaking
- composite (in particular, glueball) inflation

# Gravity Backgrounds

Solutions of 10d SUGRA equations of motion

Consistent truncation: significant simplification

(Recall: Consistent truncation means that every solution of the lower dimensional action lifts to a solution of the full 10d action)

We will investigate a 5d consistent truncation of type IIB, established in [M. Berg, M. Haack, W. Muck, hep-th/0507285]

(This encompasses MN, KS solutions, but we will look for nonsusy ones.)

→ 5d fields:

– metric:  $g_{IJ}(x^I)$

– 6 scalars:  $\phi(x^I), p(x^I), q(x^I), u(x^I), v(x^I), b(x^I)$

## 5d action:

Let us denote  $\{\varphi^i\} = \{\phi, p, q, u, v, b\}$ :

$$S = \int d^5x \sqrt{-\det g} \left[ -\frac{R}{4} + \frac{1}{2} G_{ij}(\varphi) \partial_I \varphi^i \partial^I \varphi^j + V(\varphi) \right],$$

$G_{ij}(\varphi)$  - sigma model metric ,

$V(\varphi)$  - **complicated** potential

## Equations of motion:

$$\nabla_{5d}^2 \varphi^i + \mathcal{G}^i_{jk} g^{IJ} (\partial_I \varphi^j) (\partial_J \varphi^k) - V^i = 0 ,$$

$$-R_{IJ} + 2G_{ij} (\partial_I \varphi^i) (\partial_J \varphi^j) + \frac{4}{3} g_{IJ} V = 0 ,$$

$\mathcal{G}^i_{jk}$  - Christoffel symbols for  $G_{ij}$  ,  $V^i = G^{ij} \partial_{\varphi^j} V$  .

# dS and Inflationary Solutions

Want to find a solution with the 5d metric:

$$ds_{5d}^2 = e^{2A(r)} \left[ -dt^2 + a(t)^2 d\vec{x}^2 \right] + dr^2$$

[K. Ghoroku, M. Ishihara, A. Nakamura, hep-th/0609152: Used a 10d solution in IIB with such external 5d metric and  $a(t) = e^{\sqrt{\frac{\Lambda}{3}}t}$  to study gauge theory in dS space. But the two scalars in that solution:  $\phi(r), C(r) \Rightarrow$  **not compatible with above consistent truncation.**]

Hubble parameter:  $H(t) \equiv \frac{\dot{a}(t)}{a(t)} \quad \left( \Rightarrow \dot{H} = \frac{\ddot{a}}{a} - H^2 \right)$

**Note:** • **dS space:**  $H = \text{const}$

• **Slow roll inflation:**  $H = H(t)$ , but  $\dot{H}$  small

[More precisely:  $\ddot{a} > 0 \Leftrightarrow \epsilon \equiv -\frac{\dot{H}}{H^2} < 1$ ; slow roll:  $\epsilon \ll 1$ ]



## Solving the coupled system of EoMs:

- Subtruncation of the consistent truncation:

Can consistently set  $u \equiv 0, v \equiv 0, b \equiv 0$

[EoMs identically solved]

→ Study class of solutions with **only nontrivial scalars**:

$$\phi(x^I), p(x^I), q(x^I)$$

- Look for quasi-de Sitter solutions (i.e. with  $H \approx \text{const}$ ):

In gauge/gravity duality context: these **scalar fields - glueballs**

Discrete mass spectrum → inflaton mass dynamically fixed

⇒ **No  $\eta$  problem!**

BUT:

Number of EoMs for scalar fields and metric functions  
is with one more than number of unknown functions

→ No solution?

Fortunately, **we showed:** [as long as  $A'(r) \neq 0$ ]

**One equation is dependent on the others!**

Solutions with  $H = \text{const}$ :

- 3-parameter family with  $q = -6p$  and  $\phi = 0$   
[analytical solution]
- two 4-parameter families with  $q = -\frac{3}{2}p$  and  $\phi = 3p$   
[numerical solutions]

## Solutions with time-dependent $H$ :

Look for small time-dependent deviations from exact  $H = \text{const}$  (i.e. pure  $dS$ ) solution: [ Recall:  $-\frac{\dot{H}}{H^2} \ll 1$  ]

→ Deform 5d metric ansatz and ansatz for scalar field(s)

$$\text{Metric: } ds_5^2 = e^{2A} \left[ -dt^2 + e^{2\tilde{H}} d\vec{x}^2 \right] + dr^2$$

- Expand in  $\gamma \ll 1$  around the analytical solution:

$$\phi = \gamma \phi_{(1)} + \gamma^3 \phi_{(3)} + \mathcal{O}(\gamma^5)$$

$$A = A_{(0)} + \gamma^2 A_{(2)} + \mathcal{O}(\gamma^4)$$

$$\tilde{H} = \tilde{H}_{(0)} + \gamma^2 \tilde{H}_{(2)} + \mathcal{O}(\gamma^4),$$

where  $A_{(0)}$  ,  $\tilde{H}_{(0)} = H_0 t$  - an. sol.

Leading order solution:

At order  $\gamma$ : Single equation for  $\phi_{(1)}$

At order  $\gamma^2$ : Coupled system for  $A_{(2)}$ ,  $\tilde{H}_{(2)}$ , incl.  $\phi_{(1)}^2$  terms

→ Obtain explicit solutions:  $H(t)$  and  $\phi(t)$

Recall that inflationary slow roll parameters:

$$\varepsilon = -\frac{\dot{H}}{H^2} \quad \text{and} \quad \eta = -\frac{\ddot{\phi}}{H\dot{\phi}}$$

For our solution:

$$\varepsilon \sim \mathcal{O}(\gamma^2) \ll 1 \quad , \quad \eta = 3 + \mathcal{O}(\gamma^2) \approx 3$$

## Ultra-slow roll inflation:

For inflationary parameters  $\varepsilon \ll 1$  and  $\eta = 3$ :

→ **Ultra-slow roll** inflation [arXiv:gr-qc/0503017, W. Kinney]

Gives  $n_s \approx 1$  (i.e. scale-invariant spectrum), but does not last for more than a few e-foldings

⇒ Can only be a transient phase, before usual slow roll

However, such a phase could explain **low- $l$  anomaly in CMB power spectrum**

[  $l \lesssim 40$  – power deficit ( $\sim 10\%$ ) compared to slow roll expectation ]  
( lower  $l$  – larger angular scales )

# Toward slow roll inflation:

(arXiv:1611.00295 [hep-th])

Consider higher order in  $\gamma$ :

$$\phi = \gamma \phi_{(1)} + \gamma^3 \phi_{(3)} + \mathcal{O}(\gamma^5)$$

$$A = A_{(0)} + \gamma^2 A_{(2)} + \gamma^4 A_{(4)} + \mathcal{O}(\gamma^6)$$

$$\tilde{H} = \tilde{H}_{(0)} + \gamma^2 \tilde{H}_{(2)} + \gamma^4 \tilde{H}_{(4)} + \mathcal{O}(\gamma^6)$$

Take:  $\phi_{(1)} = 0$  ,  $A_{(2)} = 0$  ,  $H_{(2)} = 0$

(Satisfies the EoMs to order  $\gamma^2$ .)

So to leading order in  $\gamma$ :  $\phi \sim \mathcal{O}(\gamma^3) \rightarrow \phi^2 \sim \mathcal{O}(\gamma^6)$

$\Rightarrow$  EoMs for  $A_{(4)}$ ,  $\tilde{H}_{(4)}$  decouple from  $\phi_{(3)}$  !

(Ultra-slow roll solution above:  $A_{(2)}$ ,  $\tilde{H}_{(2)}$ ,  $\phi_{(1)}^2 \sim \mathcal{O}(\gamma^2)$ .)

## Toward slow roll inflation:

→ Solutions for  $A_{(4)}$ ,  $\tilde{H}_{(4)}$  are independent of  $\phi_{(3)}$  and vice versa

In particular, integration constants in  $\phi_{(3)}$  are unrelated to those in  $A_{(4)}$ ,  $\tilde{H}_{(4)}$

This relaxes constraint that led to ultra-slow roll at  $\mathcal{O}(\gamma^2)$

⇒ Can show that, for suitable choices of the integration constants, one can obtain:

$$\eta \ll 1$$

→ Slow roll regime !

# Summary

Found so far:

- Three multi-parameter solutions in 5d consistent truncation of IIB supergravity

[ $dS_4$  space fibered over the fifth dimension]

- Ultra-slow roll glueball inflation model [ $t$ -dep. deformation]
- Showed that it is possible to obtain slow roll regime

Open issues:

- Slow roll Glueball Inflation ?...
- Microscopic realization ?...
- Inflaton mass (mass-spectrum of fluctuations) ?...



**Thank you!**